

Radiation Pressure and Turbulence in AGN Tori

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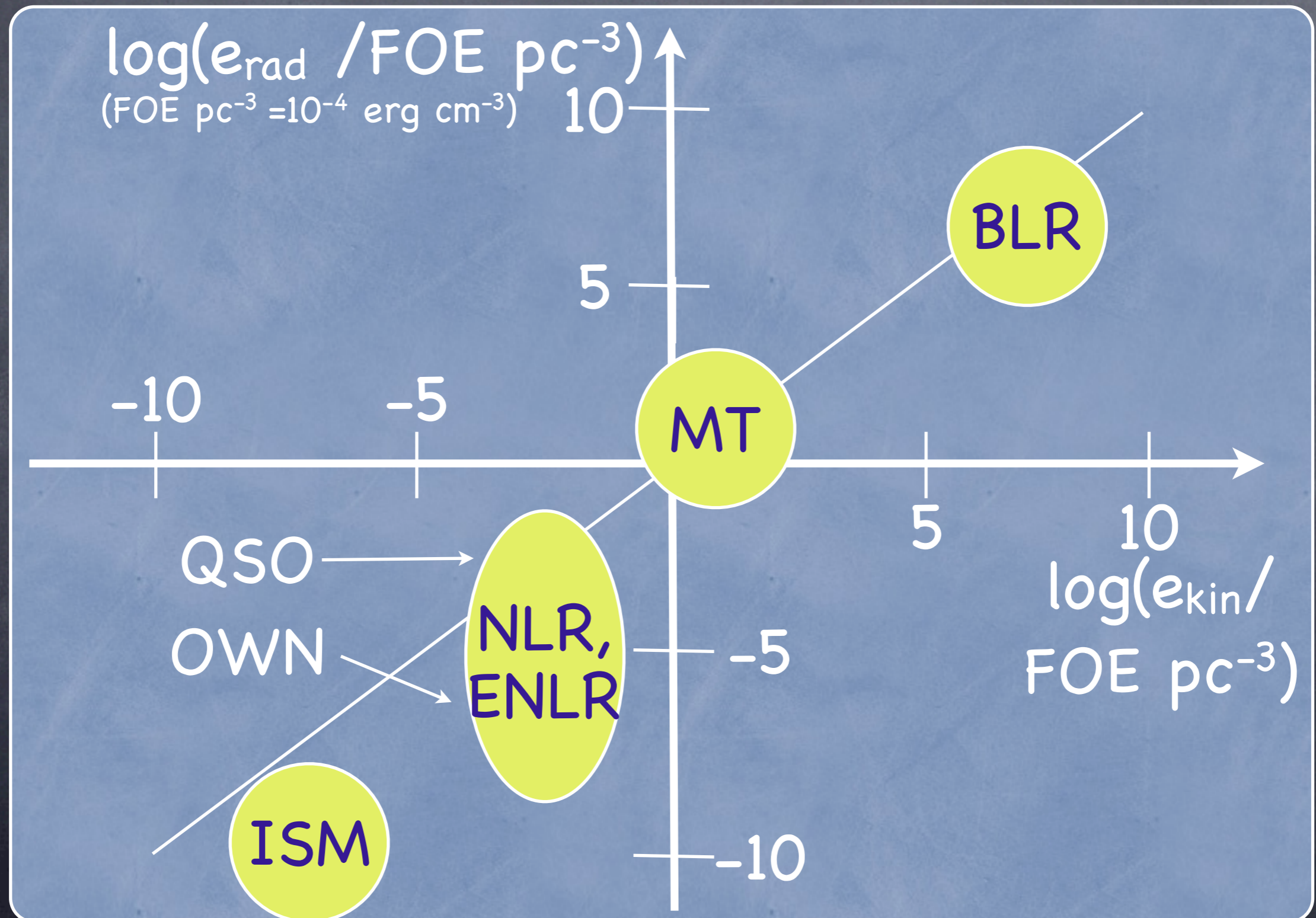
with Marc Schartmann, Andreas Burkert
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Physics of Galactic Nuclei - Ringberg, June 2009

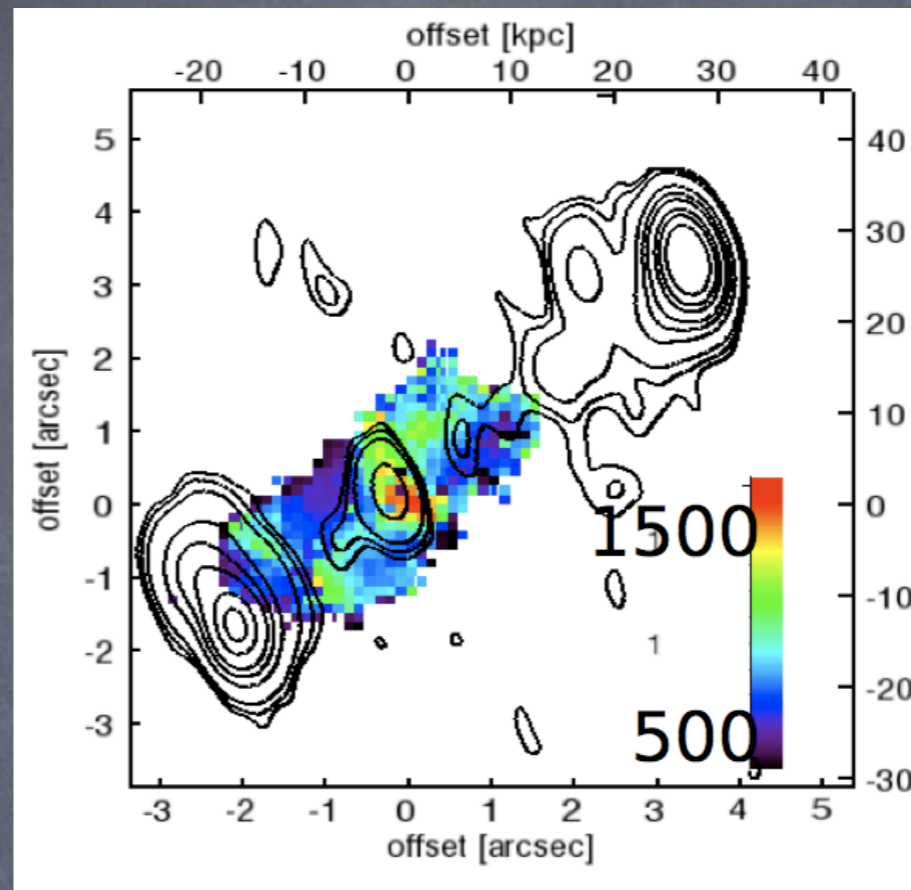
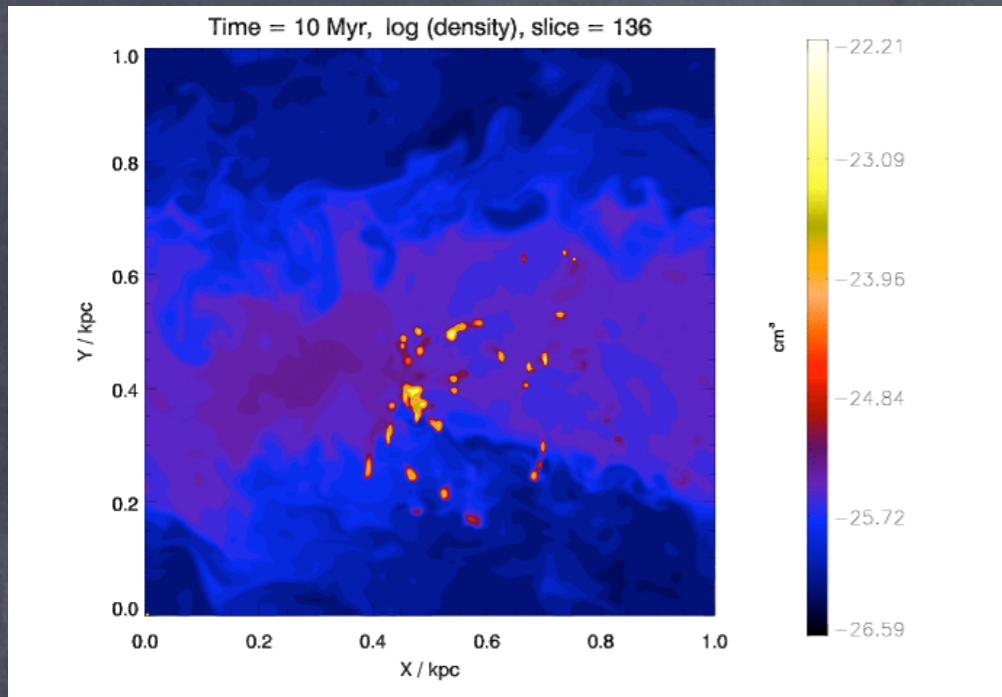
Overview

- Radiative and kinetic energy in the ISM
- The literature
 - the geometrical thickness problem
 - the importance of radiation pressure
- A toy model for radiative interactions
 - enhanced elasticity
 - radiatively driven turbulence

Multi-phase turbulence: radiative versus kinetic



Extended Narrow Line Regions



0406-242:
 $z=2.4$,
[OIII], $H\beta$,
 $>10^{10} M_{\odot}$,
 $v \approx 1000$ km/s
(Nesvadba et al. 2008)

Simulation:

- hydro (grid) + o.thin cooling
- contains all three phases (rel., hot, cold)
- e.g. Krause & Alexander 2007

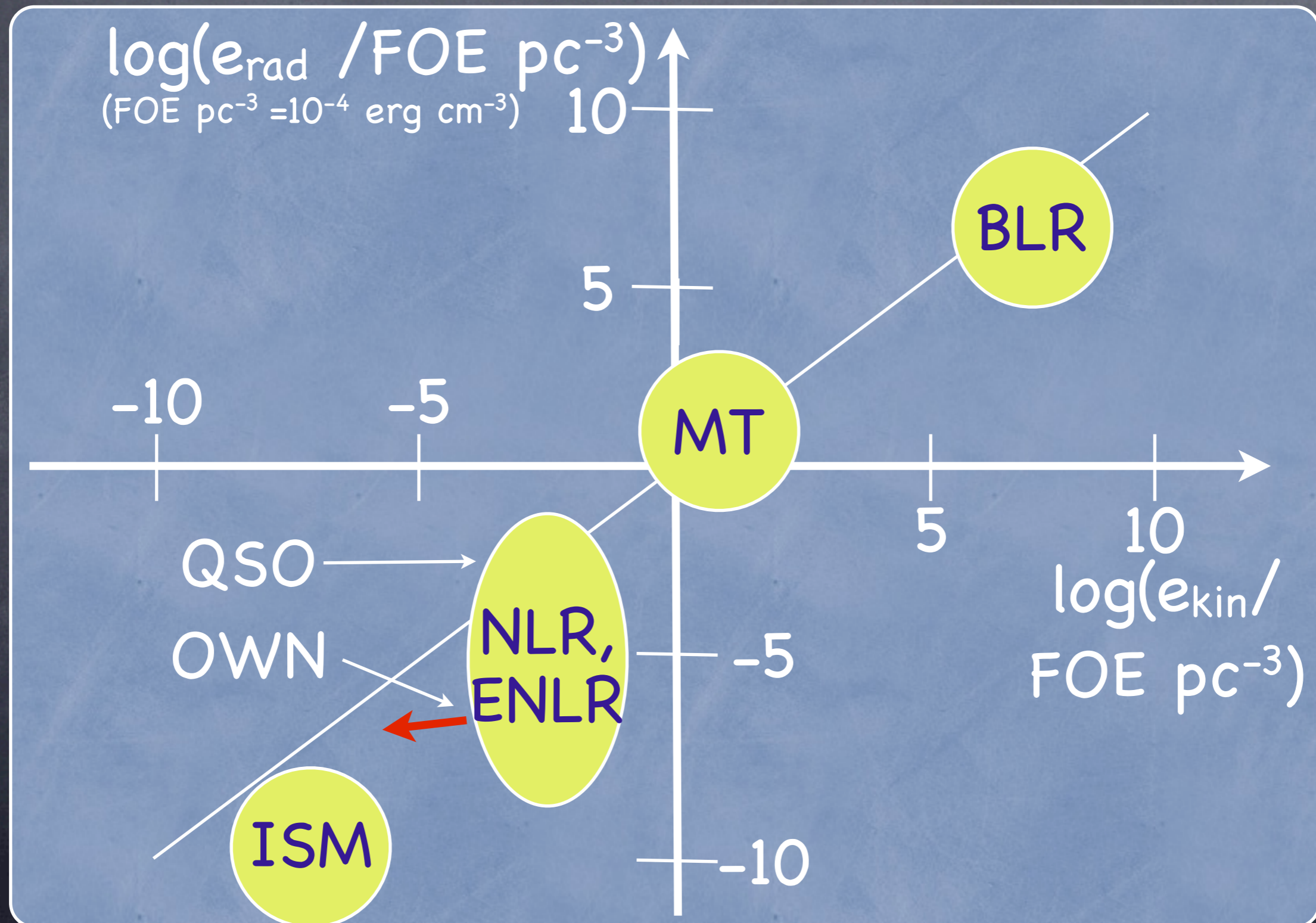
Results:

- EL-gas survives stirring
- Phase equilibrium depends on energy in box
- Decay due to shocks/radiation

- massive star forming galaxies at $z = 2-5$
- gas expelled permanently by jet feedback
- Radiative decay, 100 Myr

Evolves towards equilibrium, i.e. produces radiation by reducing turbulent energy

Multi-phase turbulence: radiative versus kinetic



Geometry known: example: NGC 1068

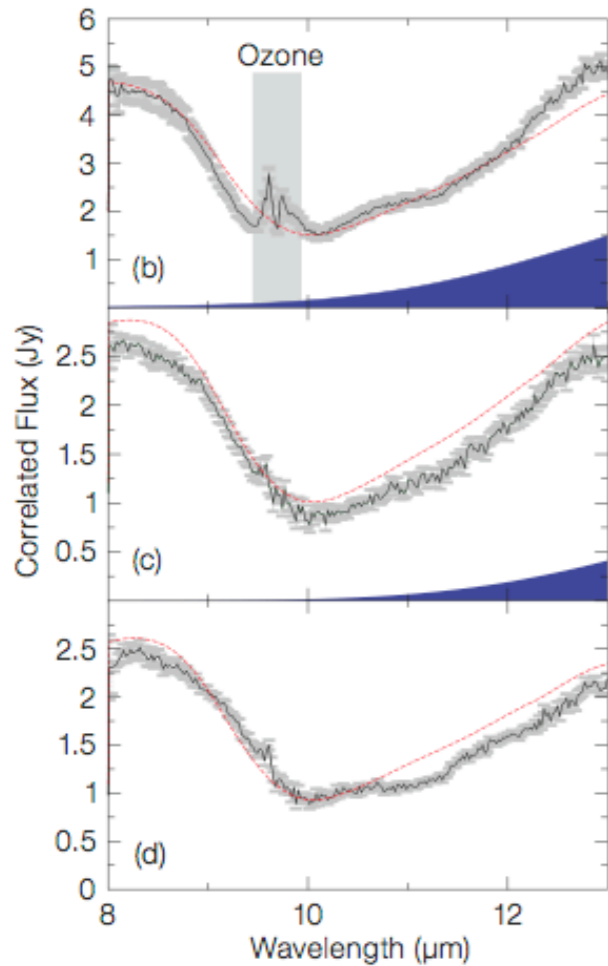
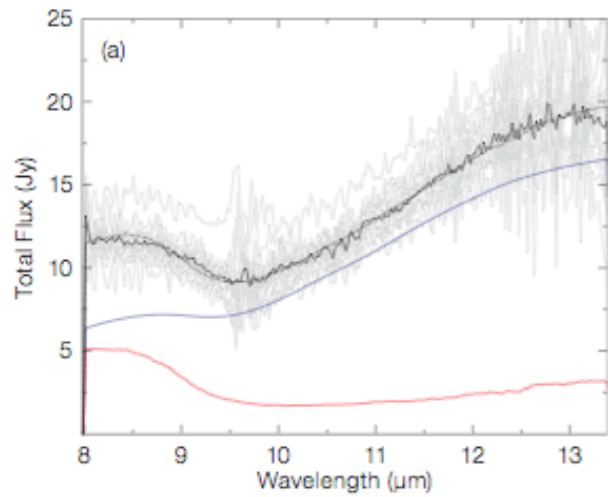


Figure 1. Results of MIDI observations of NGC 1068. (a) Total flux $F_{\text{tot}}(\lambda)$: the contribution of the hot component is shown in red, that of the extended component in blue. (b) Correlated flux $F_{\text{corr}}(\lambda)$ obtained with a 40-m baseline orientated along position angle P. A. = 36° . The red dotted line gives the model fit and the blue shaded area shows the contribution of the extended component. (c) $F_{\text{corr}}(\lambda)$ for 52 m baseline along P. A. = 112° . (d) $F_{\text{corr}}(\lambda)$ for 97 m baseline along P. A. = 36° . The comparison between (b) and (c) shows that the hot component is more extended (better resolved) in SE-NW direction.

- ① IR interferometry
- ① size: 1pc
- ① shape: torus
- ① hot interior
- ① cooler outskirts

← Meisenheimer et al. 2008 →

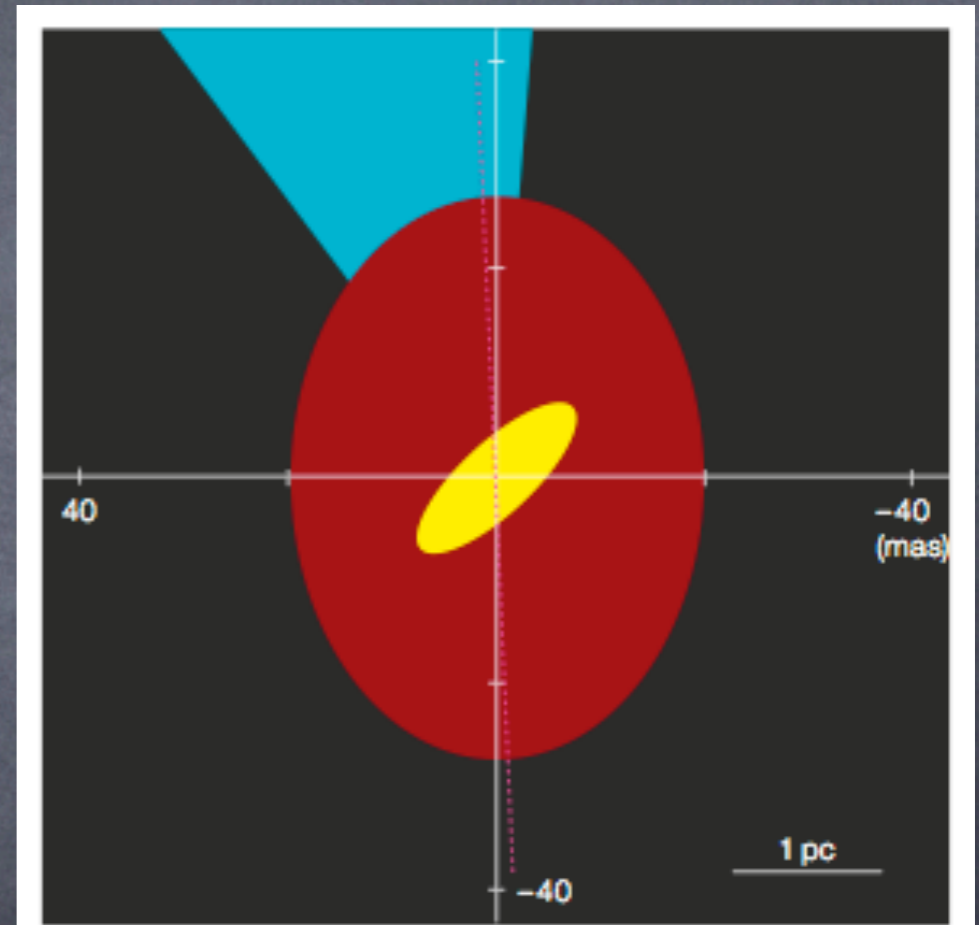


Figure 2. Observational model of the dust torus in NGC 1068. A hot component (yellow) is embedded in an extended cooler component (brown). The orientation of the radio axis is indicated by a purple dotted line and the blue wedge gives the opening angle of the ionisation cone, observed on 100-pc scales.

The literature

- Krolik & Begelman 1988
 - obs: $h/r \approx 0.7$ & dust, $\Rightarrow v \approx 100 \text{ km/s}$, if thermal \Rightarrow no dust, \Rightarrow clumpy
 - inner edge: sublimation radius
 - cloud merger & tidal shear \Rightarrow cov.fac. $> \approx 1$, size dist., $N \approx 10^{24} \text{ cm}^{-2}$ @ Jeans limit
 - accretion due to cloud-cloud collisions \Rightarrow reduces scale height \Rightarrow thin disc
 - stirring by stellar processes not enough
 - suggestion: elastic clouds, i.e there are no hard collisions, due to magnetic fields

The literature

- Beckert & Duschl 2004
 - accretion torus model, viscosity due to cloud-cloud collisions
 - relies on magnetic field model to produce sufficient vertical turbulence
- Hönig & Beckert 2007
 - Expands BD04: direct UV radiation pressure on torus clouds
 - consistent with $\theta \propto 1/L$ (receding torus model)
 - still imply magnetic clouds to isotropize motions

The literature

- Pier & Krolik 1992
 - radiation pressure dominates over central gravity in the hot dust region
 - fluctuating part of radiation force might drive turbulence
- Krolik 2007, Shi & Krolik 2008
 - magnetic cloud model: "required field strengths are not terribly plausible"
 - self-consistent solutions, IR-radiation pressure supported torus
- Krolik et al 2007, Blaes et al 2007
 - MRI turbulent disk increase scale height by radiation pressure

May the radiation force
drive turbulence or
make the clouds elastic?

- toy model -

A nice new formalism: radiative potential

$$\mathbf{F}_{ij} = \frac{L_j}{4\pi r_{ij}^2 c} \pi R_i^2 \mathbf{e}_{ij},$$

$$\tilde{m}_i = \sqrt{\frac{\pi\sigma}{cG}} T^2 R_i^2,$$

$$\mathbf{F}_{ij} = G \frac{\tilde{m}_i \tilde{m}_j}{r_{ij}^2} \mathbf{e}_{ij},$$

$$E_{\text{pot}} = G \frac{\tilde{m}_i \tilde{m}_j}{r_{ij}^2}.$$

- two clouds i,j
- force $\propto 1/r^2$
- introduce radiative mass
- symmetric material properties
- here: optically thick
- analogy: gravity

Two cloud system (uniform clouds)

$$\frac{1}{2} m v_c^2 = G \frac{\tilde{m} \tilde{m}}{R}$$

$$\frac{\tilde{m}}{m} = \frac{3}{4} \sqrt{\frac{\sigma}{\pi c G} \frac{T^2}{\rho R}}$$

$$= 2,300 \left(\frac{T}{1000 \text{ K}} \right)^2 \left(\frac{\rho}{10^{-17} \text{ g cm}^{-3}} \right)^{-1} \left(\frac{R}{10^{-3} \text{ pc}} \right)^{-1}$$

- clouds need high T => radiative > gravitational mass

- approaching clouds bounce like protons

- require $v > v_c$ for direkt hit**

$$v_c = \sqrt{\frac{3\sigma}{2c} \frac{T^2}{\rho^{1/2}}}$$

$$= 168 \text{ km s}^{-1} \left(\frac{T}{1000 \text{ K}} \right)^2 \left(\frac{\rho}{10^{-17} \text{ g cm}^{-3}} \right)^{-1/2}$$

Optically thick cloud ensemble

$$\lambda n_{\text{cl}} \pi R^2 = 1,$$

$$N = n_{\text{cl}} (4/3) \pi \lambda^3 = (4/3)^3 f^{-2},$$

$$E_{\text{pot}}(r) = \int_0^{\infty} G \frac{\bar{\rho}(r') 4\pi r'^2 dr'}{(r - r')^2}.$$

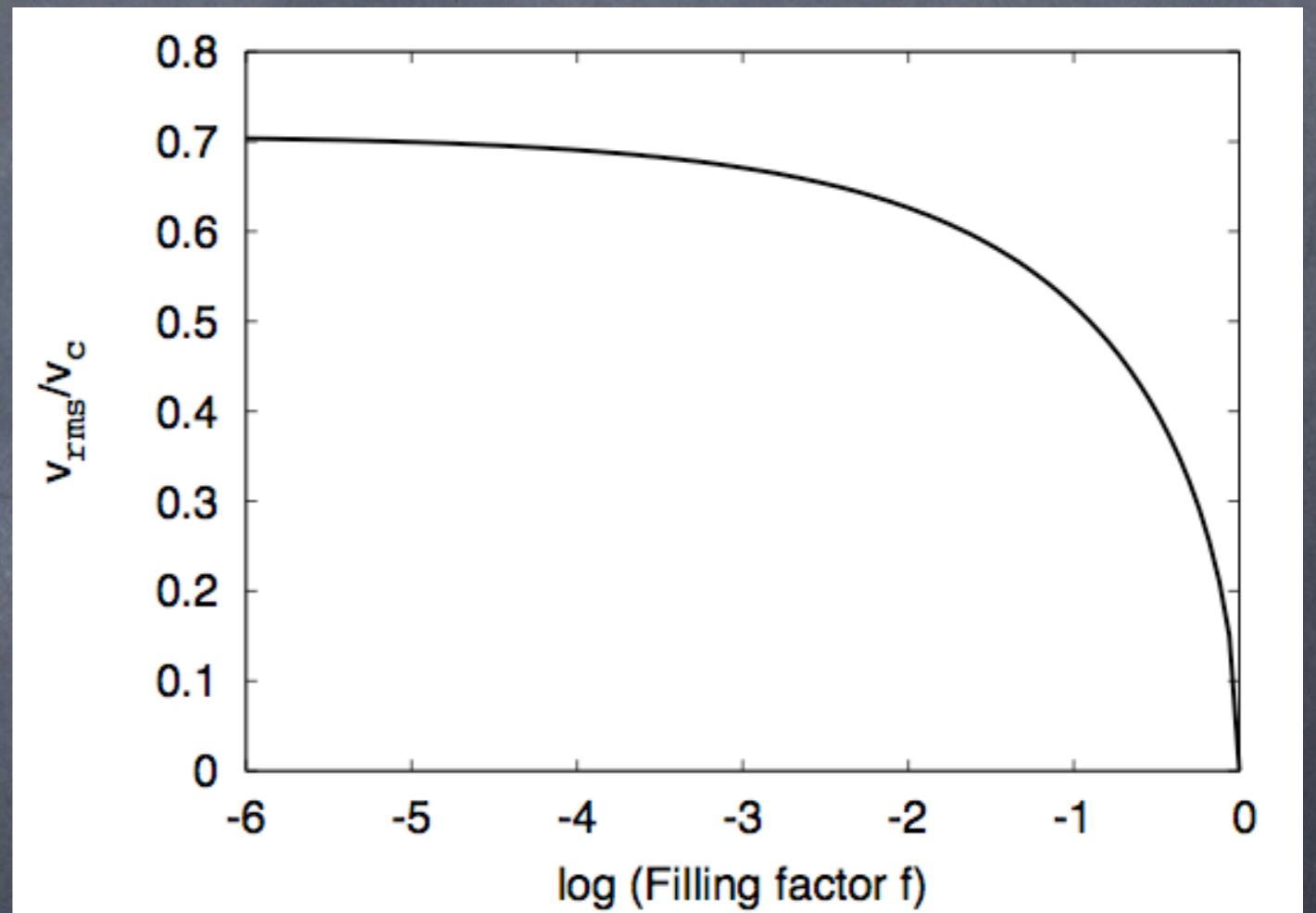
- Define horizon λ (=mean free path)
- filling factor f small (clumpy) \Rightarrow number of clouds within λ is large
- all clouds within sphere (λ) interact with origin
- may define radiative mass density
- From shell potential, force within smooth shells = 0

Optically thick cloud ensemble

- Two body interactions dominate
- Shadow effects negligible =>> Radiation force is conservative
- Virial theorem: $2T = \Delta E$

• =>>>

$$v_{\text{rms}} = \sqrt{\frac{1 - f^{1/3}}{2}} v_c$$



$$\Delta E = G\tilde{m}^2 \left(\frac{1}{R} - \frac{1}{\bar{d}} \right) = G\frac{\tilde{m}^2}{R} (1 - f^{1/3})$$

Other cloud types

- Easily find cloud parameters where radiation dominates over cloud gravity
- All examined cloud models in principle permit rms-velocities > 100 km/s @ $T(\text{dust}) = 1000\text{K}$
- Opt. thin dusty clouds unstable against own radiation pressure @ $T > 30\text{K}$

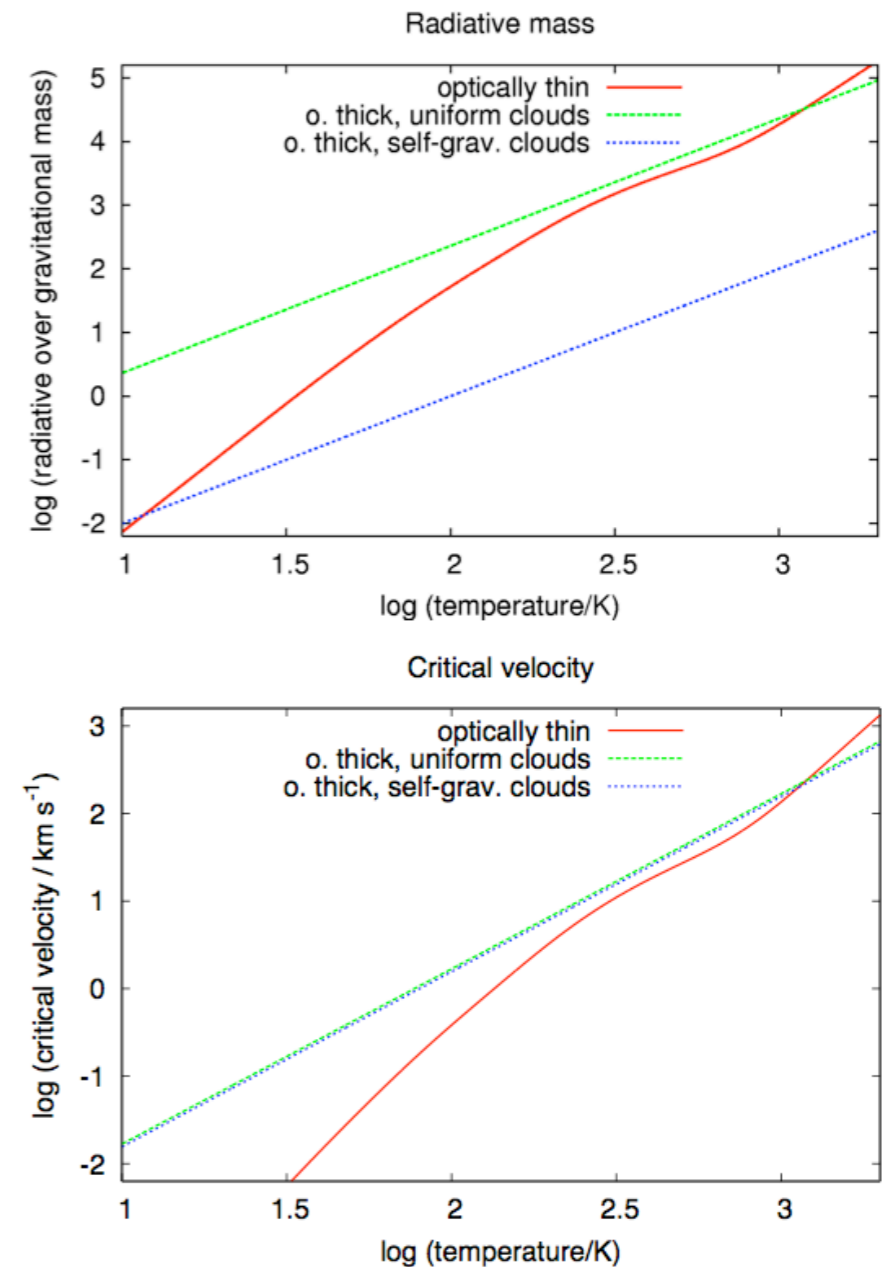


Figure 2. Radiative mass and critical velocity for reference clouds (uniform optically thick: $R = 10^{-3}$ pc, $\rho = 10^{-17}$ g cm⁻³, self-gravitating optically thick: $R = 2 \times 10^{-2}$ pc, $c_s = 1$ km s⁻¹, optically thin: $R = 10^{-2}$ pc, $\rho = 10^{-21}$ g cm⁻³).

Conclusions

- Rarely find $E_{\text{rad}} > E_{\text{kin}}$
- In torus & BLR: $E_{\text{rad}} \approx E_{\text{kin}}$
- In BLR radiation pressure thought to be important
- Presented toy model of fluctuating radiation force
- For reasonable cloud parameters, find

- enhanced elasticity
- turbulence $> 100 \text{ km/s}$

$a_{\text{cloud}} = 0.001 \text{ pc}$
 $\rho = 3 \cdot 10^{-17} \text{ g/cm}^3$
 $T_{\text{sub}} = 1000 \text{ K}$
 $M_{\text{bh}} = 7.9 \cdot 10^6 M_{\text{sun}}$
 $\text{eddratio} = 0.005$
 $n_{\text{clouds}} = 5000$

