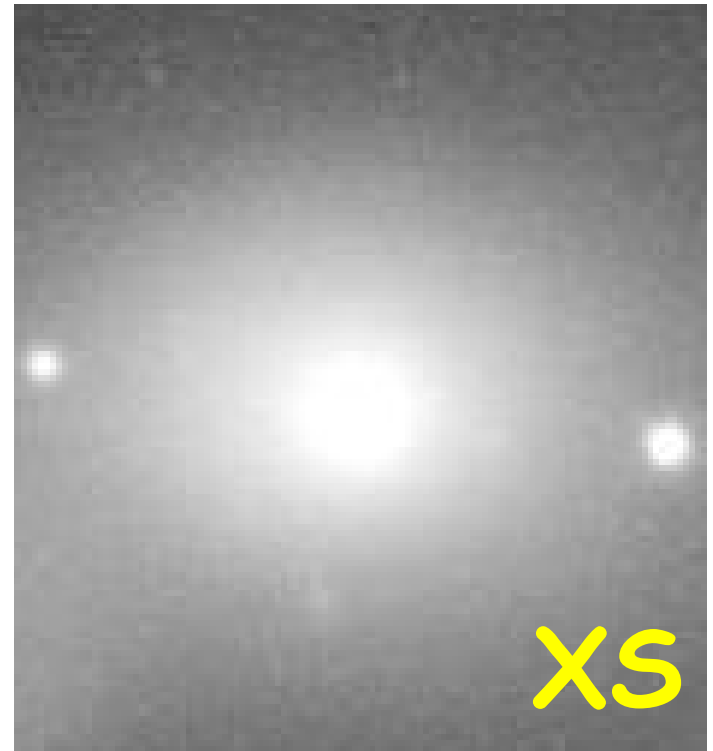
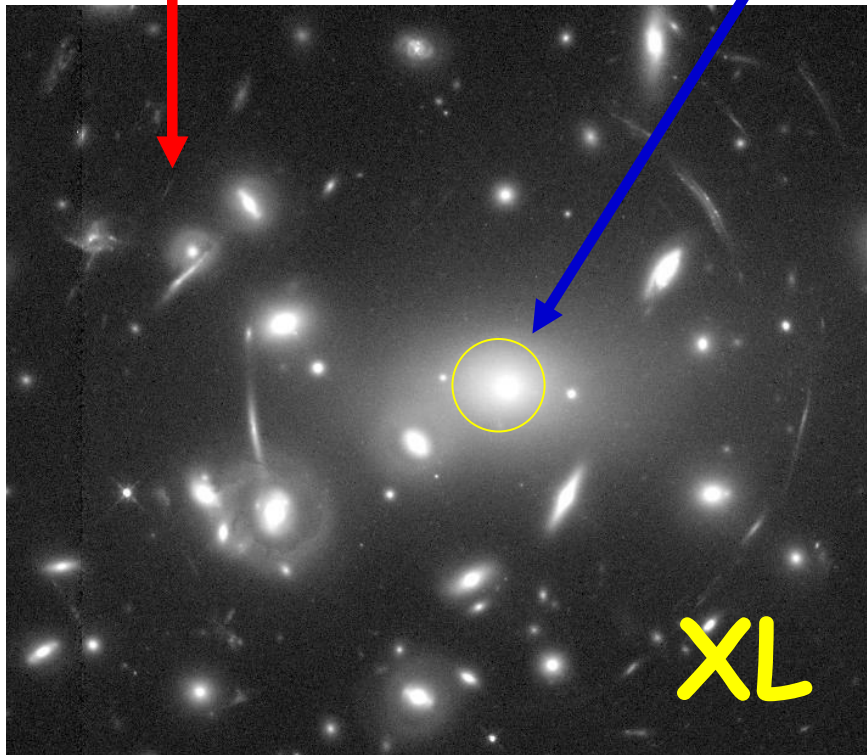


## Extra physics (cosmic rays et al.)

E.Churazov, W.Forman, A.Vikhlinin, S.Tremaine,  
O.Gerhard, C.Jones MNRAS, 2008, v388, p1062

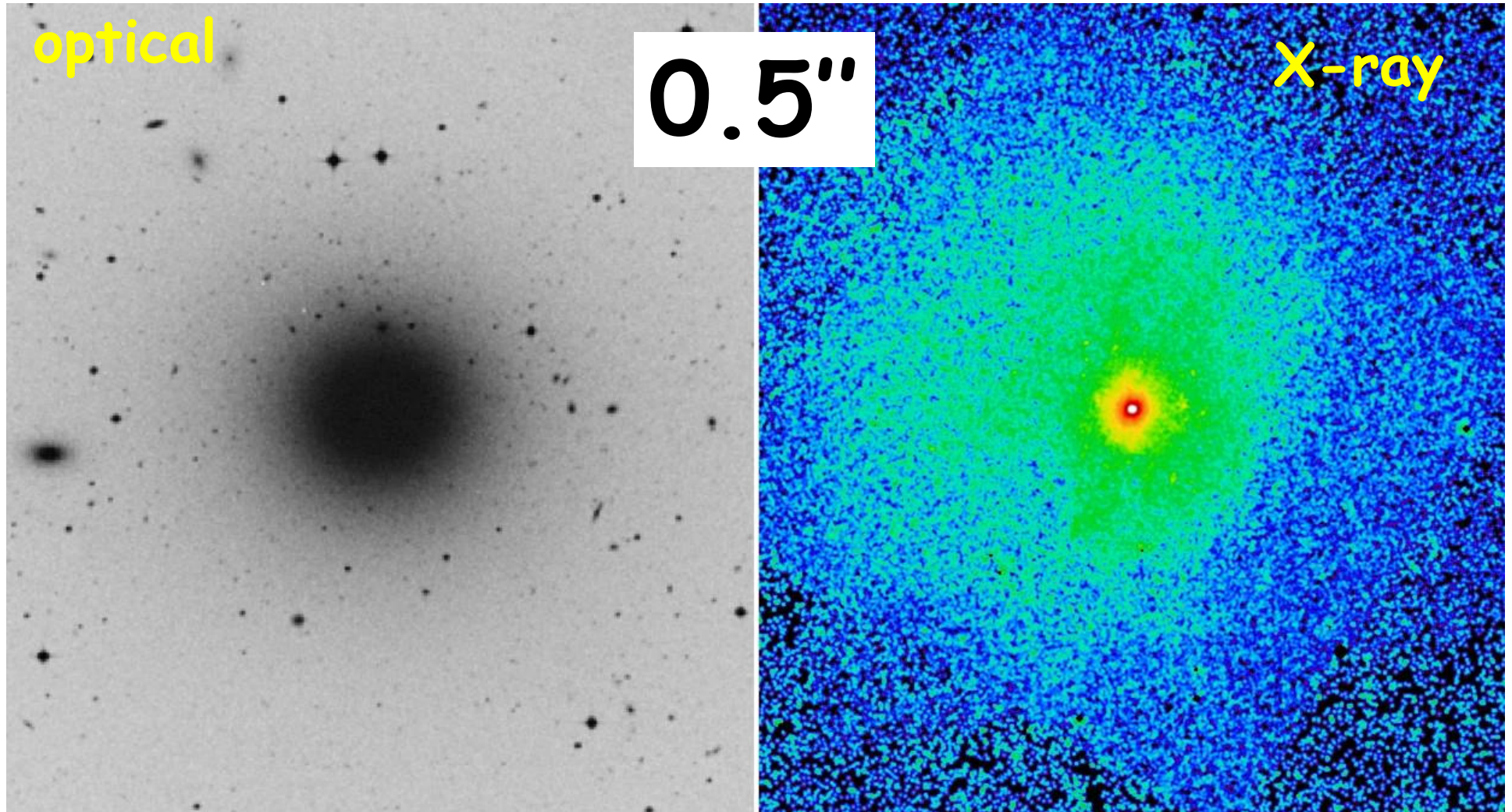
This conference

This talk



Goal: not the mass, but constraints on non-thermal pressure and on the deviations from hydrostatic equilibrium

# NGC1399 - Fornax cluster



Stars instead of lensing:  
gravity only

Gas: gravity, magnetic fields, cosmic  
rays, turbulent motions

Stars: Jeans equation [stationary, spherical system with isotropic velocity dispersion]

Gas: hydrostatic equilibrium

$$\frac{1}{n_*} \frac{dn_* \sigma^2}{dr} = - \frac{GM}{r^2}$$

$$\frac{1}{\rho_{gas}} \frac{dP}{dr} = - \frac{GM}{r^2}$$

$$P_{thermal} = nkT \quad \text{Gas thermal pressure: can be easily measured}$$

$$P = P_{thermal} + \underbrace{P_{CR} + \frac{B^2}{8\pi} + P_{turb}}_{\text{Non-thermal pressure}}$$

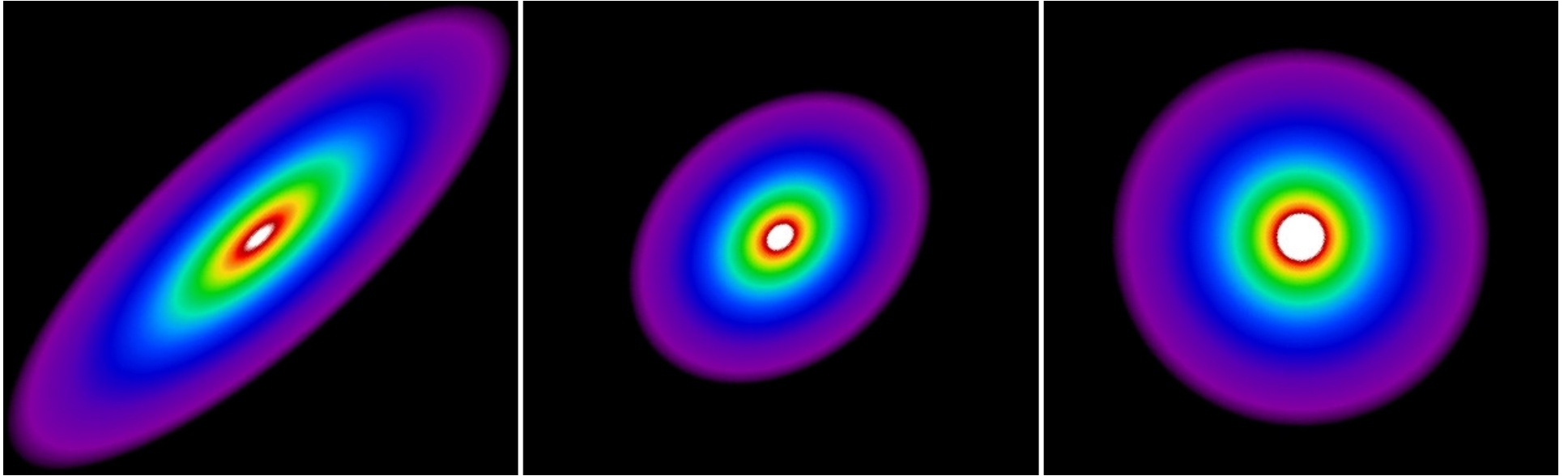
Non-thermal pressure

**Goal: to place constraints on the non-thermal pressure and on deviations from the hydrostatic equilibrium**

- 1. (X-ray) Deprojection of non-spherical clusters**
- 2. Using potential instead of mass**
- 3. Impact of non-thermal pressure**
- 4. Results and conclusions**



## Triaxial $\beta$ -model, isothermal gas, relaxed

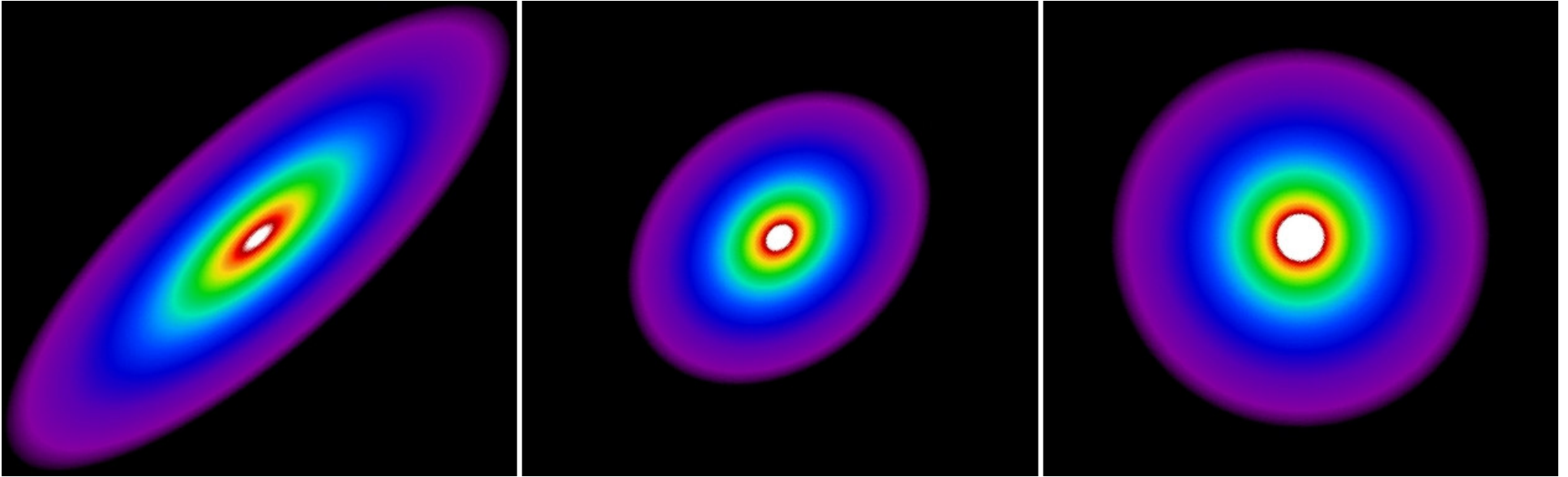


$$n_e = \left[ 1 + \left( \frac{x}{a} \right)^2 + \left( \frac{y}{b} \right)^2 + \left( \frac{z}{c} \right)^2 \right]^{-\frac{3}{2}\beta} \quad r \gg a, b, c \quad n_e = f(\theta, \phi) r^{-\gamma}$$

$$n_e \propto e^{-\frac{\mu m_p \phi}{kT}} \Rightarrow \phi = \frac{kT}{\mu m_p} [\gamma \ln r - \ln f(\theta, \phi)]$$

$$\rho = \frac{kT}{4\pi G \mu m_p} \left[ \frac{\gamma}{r^2} - \nabla^2 \ln f(\theta, \phi) \right]$$

## Triaxial $\beta$ -model, isothermal gas



$$M(< R) = \int_{|r|<R} d\vec{r} \rho = \frac{kT}{\mu m_p G} \left[ \gamma R - \int_{|r|<R} d\vec{r} \nabla^2 \ln f(\theta, \phi) \right] = \frac{kT}{\mu m_p G} \gamma R$$

From X-ray analysis:  $n_e \propto f(\theta, \phi) r^{-\gamma} \Rightarrow n_e \propto r^{-\gamma}$ ,  $T$

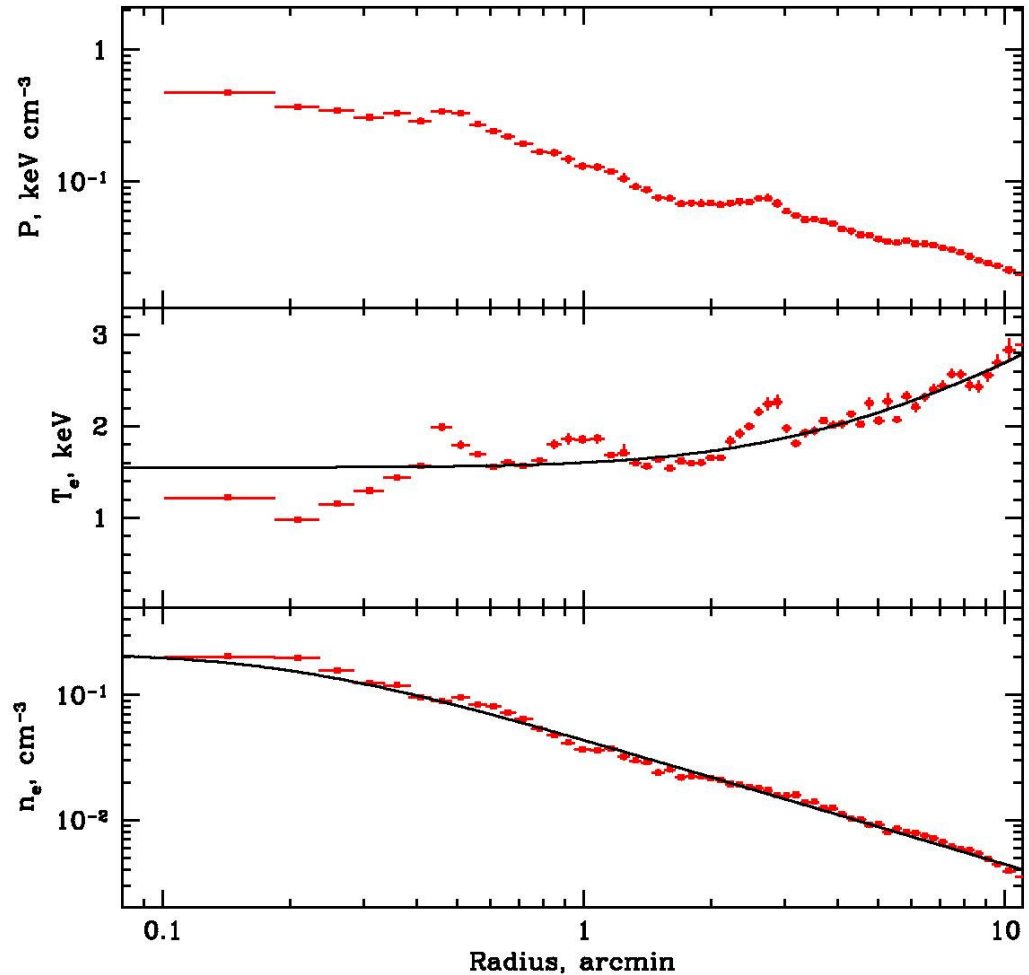
$$M_x(< R) = -\frac{R^2}{G} \frac{1}{\mu m_p n_e} \frac{dP}{dr} = \frac{kT}{\mu m_p G} \gamma R$$

$T$ ;  $n_e \propto f(\theta, \phi) g(r) \propto f(\theta, \phi) r^{-\gamma}$  e.g.  $\beta_3$ -model, large  $r$

# Calculate potential rather than mass!

$$\frac{GM}{r^2} = -\frac{1}{\rho} \frac{dP}{dr}$$

$$\frac{d\phi}{dr} = -\frac{1}{\rho} \frac{dP}{dr}$$



$$\phi = -\frac{k}{\mu m_p} \left[ \int T_e \frac{d \ln n_e}{dr} dr + T_e \right] + C \quad \phi = -\frac{kT_e}{\mu m_p} \ln n_e + C$$

$$\frac{d\varphi_x}{dr} = - \frac{1}{\mu m_p n} \frac{dP}{dr}$$

**We measure** :  $n, T, P = nkT$

**What can be wrong?:**  $n, T, P = nkT + ..$





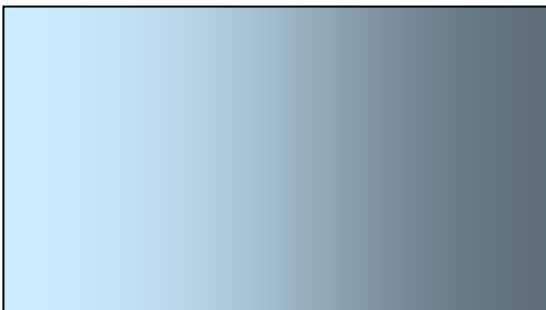
Uniform pure thermal gas  
 $\rho$ -OK,  $P$ -OK



Uniform mixture: thermal + CRs  
 $\rho$ -OK,  $P$ -wrong

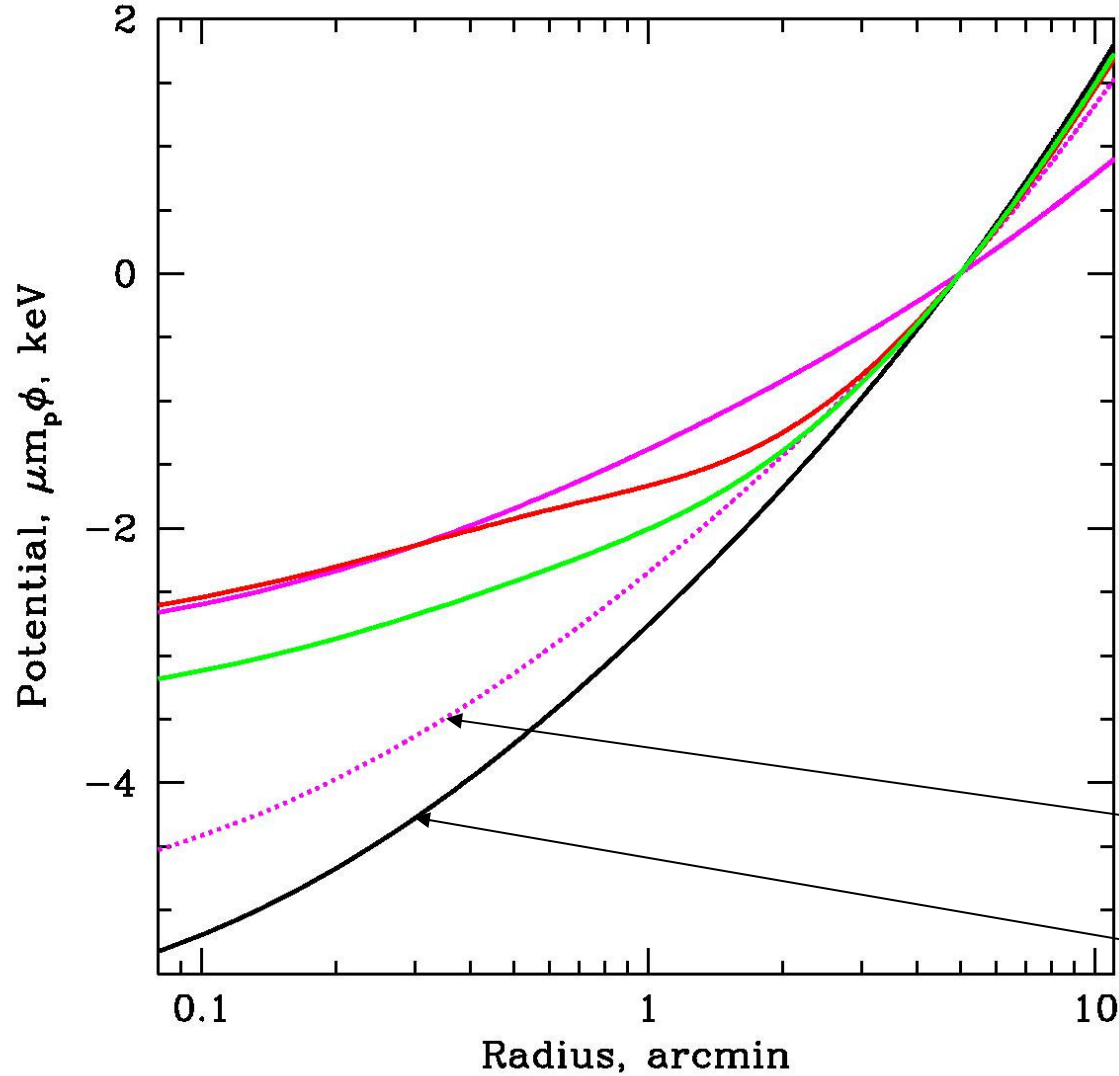


Bubbles of CRs in thermal gas  
 $\rho$ -wrong,  $P$ -wrong



Fraction of CRs varies with radius  
 $\rho$ -wrong,  $P$ -wrong

# Impact of non-thermal pressure



$$\frac{d\varphi_{true}}{dr} = -\frac{1}{\rho_{true}} \frac{dP_{true}}{dr}$$

$$P_{thermal} = \alpha P_{true}$$

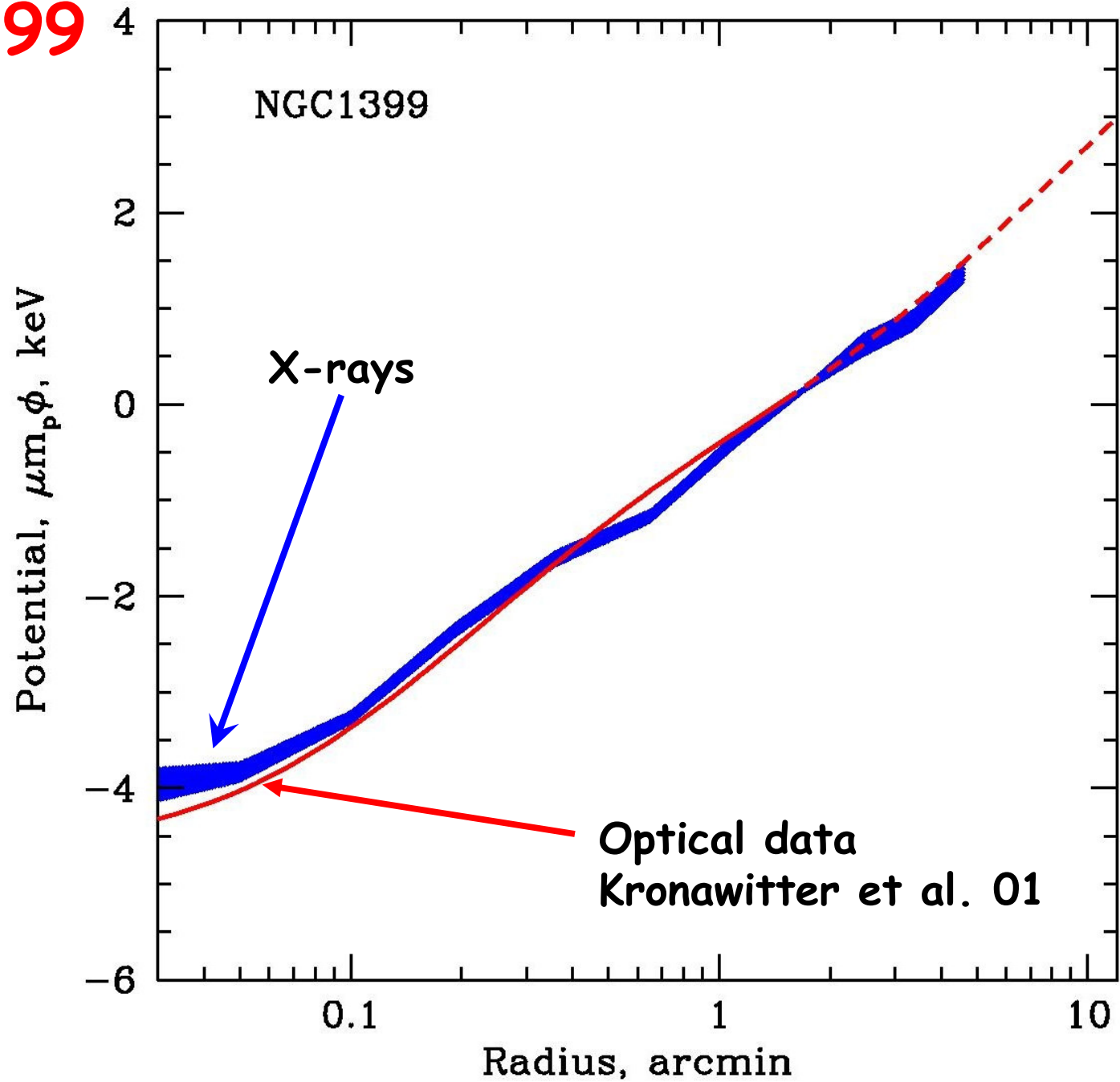
$$\frac{d\varphi_X}{dr} = -\frac{1}{\rho_{true}} \frac{d\alpha P_{true}}{dr}$$

"Observed" potential  
25% - cosmic rays

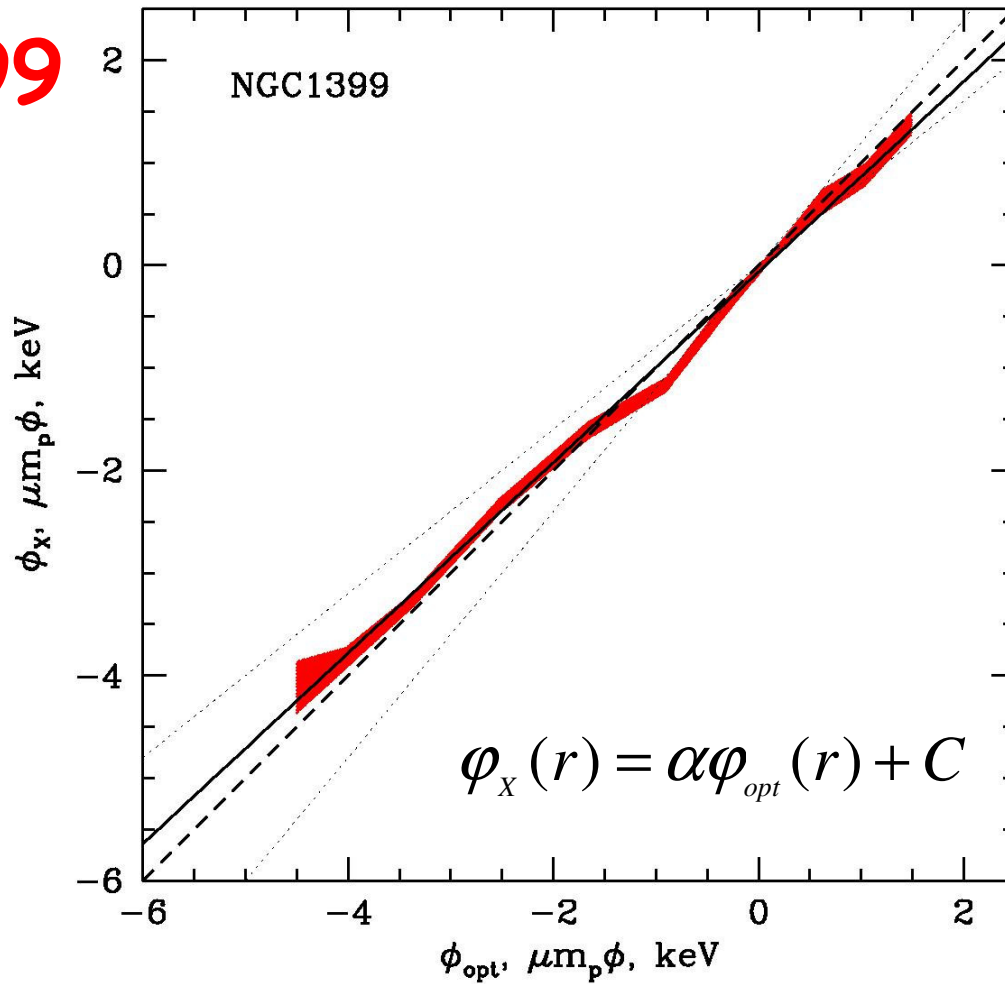
True potential

$$\varphi_X(r) \approx \alpha \varphi_{true}(r), \quad \alpha \leq 1$$

# NGC1399

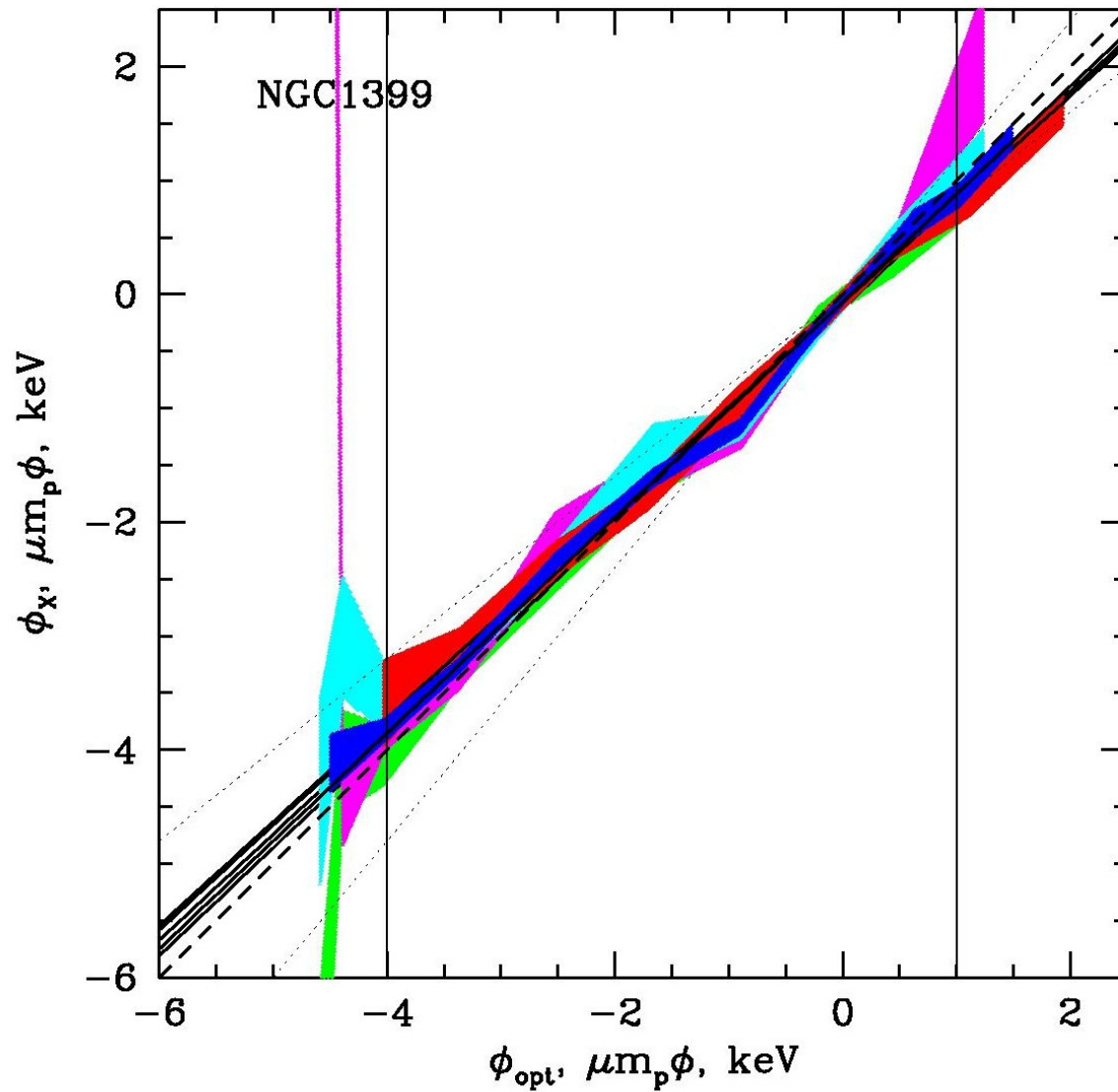


NGC1399



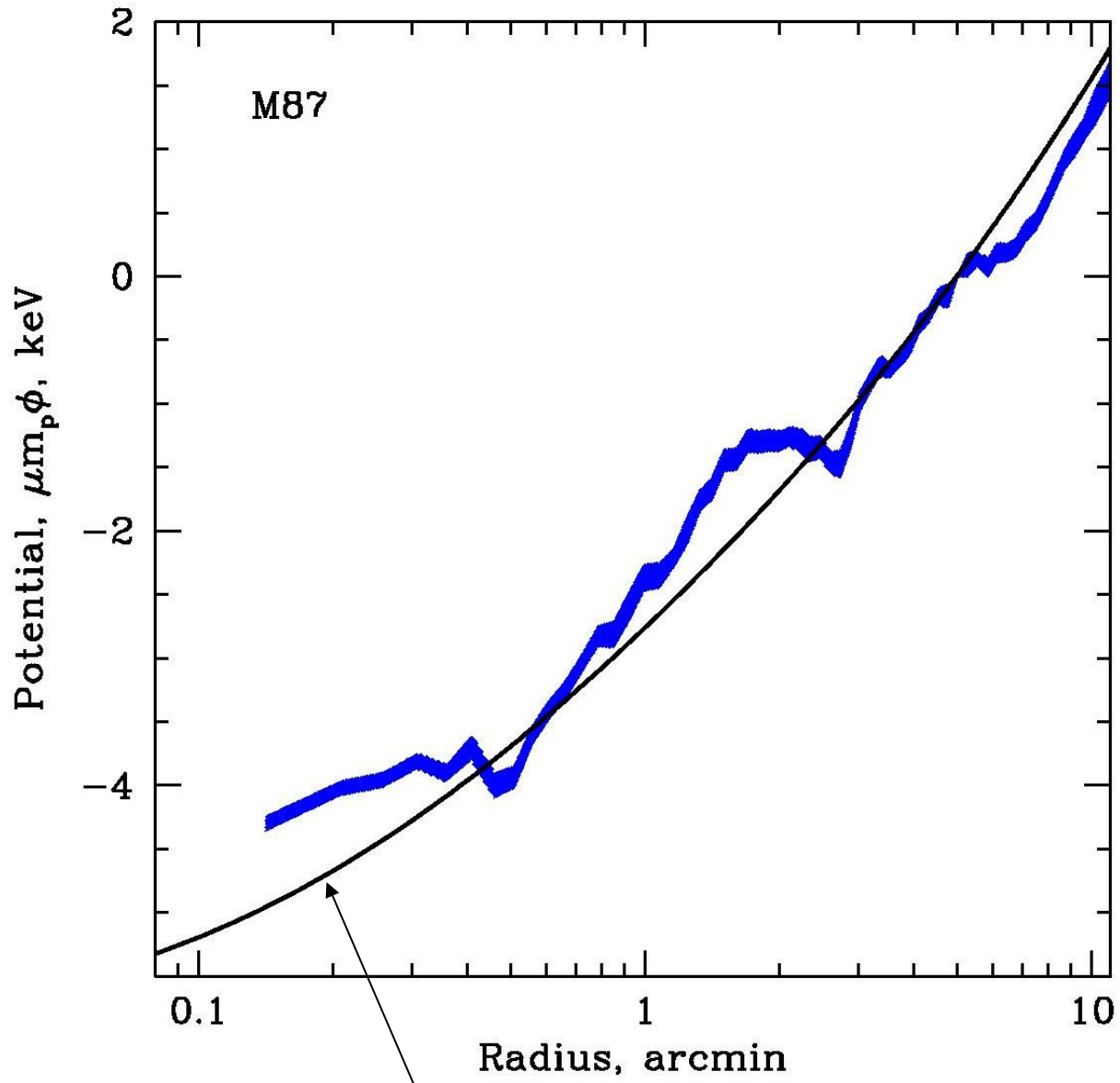
$$\varphi_X(r) \approx 0.93 \varphi_{opt}(r) + C$$

$$U_{CR} + \frac{H^2}{8\pi} + U_{turb} = 0.07 U_{thermal}$$



Four  $90^\circ$  wedges={0.95, 0.93, 0.92, 0.95}, averaged over  $360^\circ$ : 0.93

**M87**

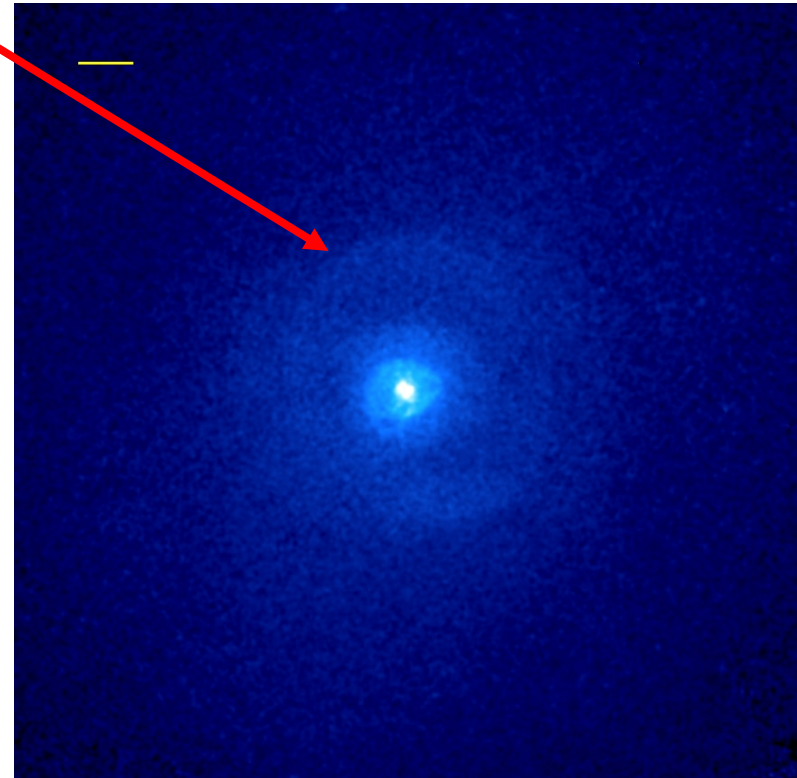


Potential from stellar kinematics [Romanowsky & Kochanek, 2001]



# Shock wave in M87 (Forman et al., 2007)

X-ray observations

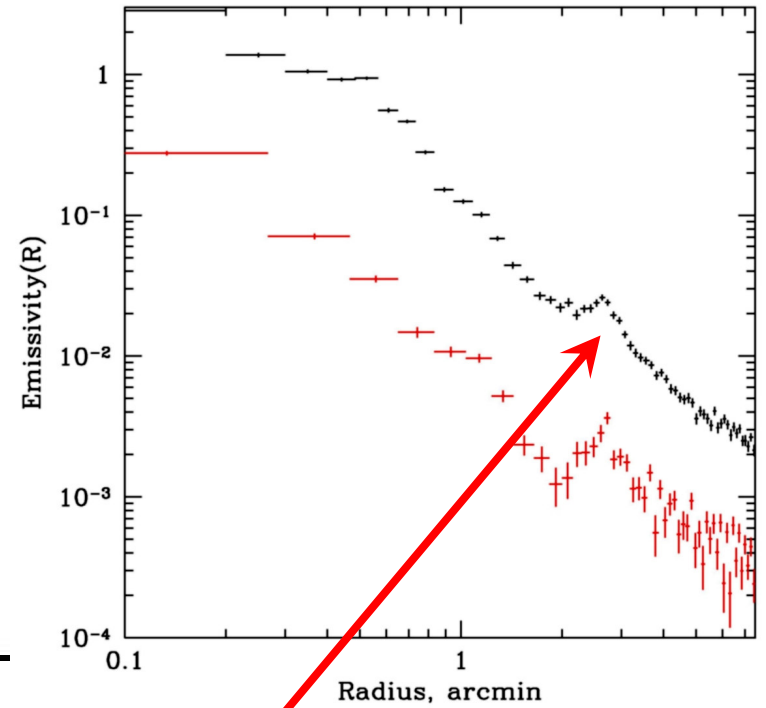
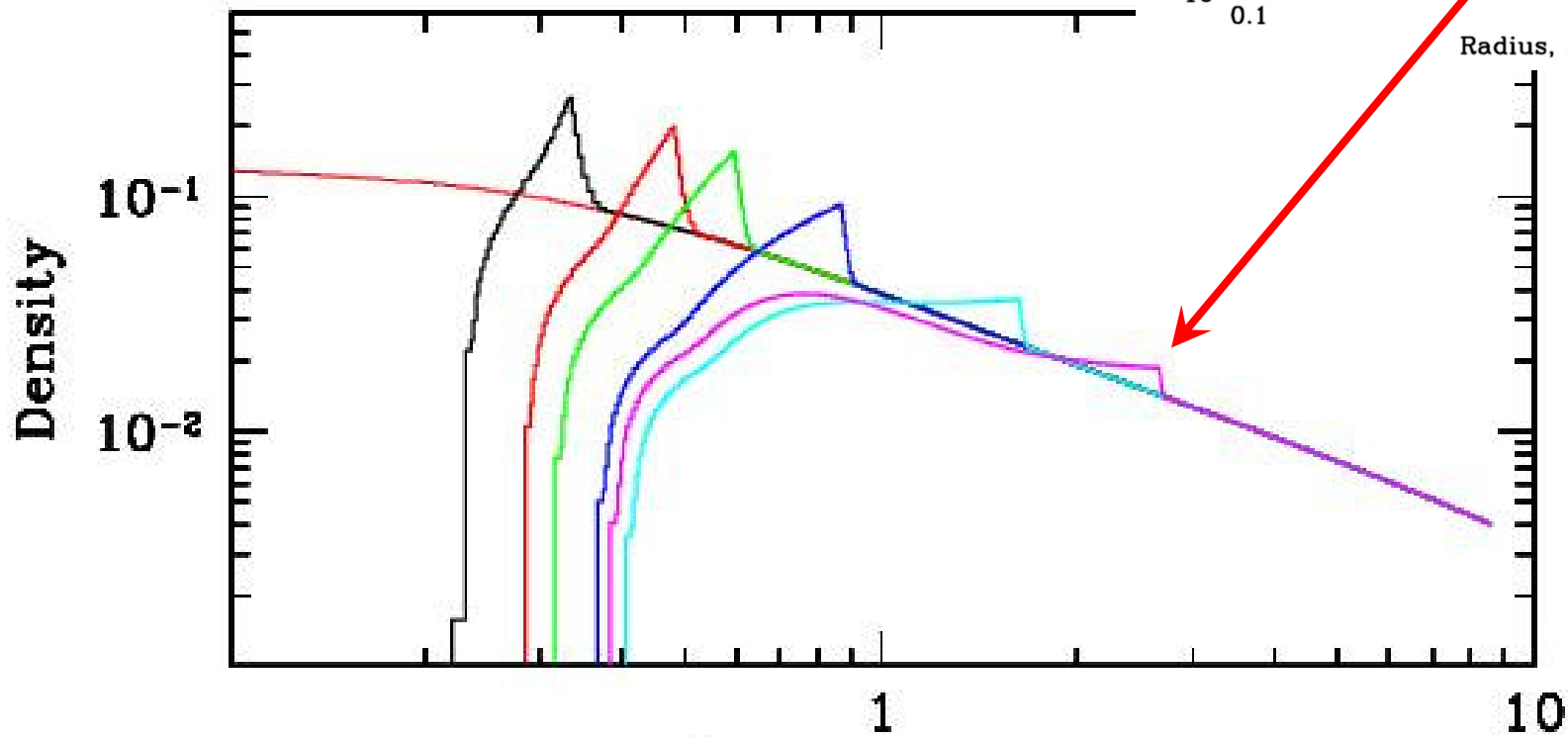


Ruszkowski+, 04

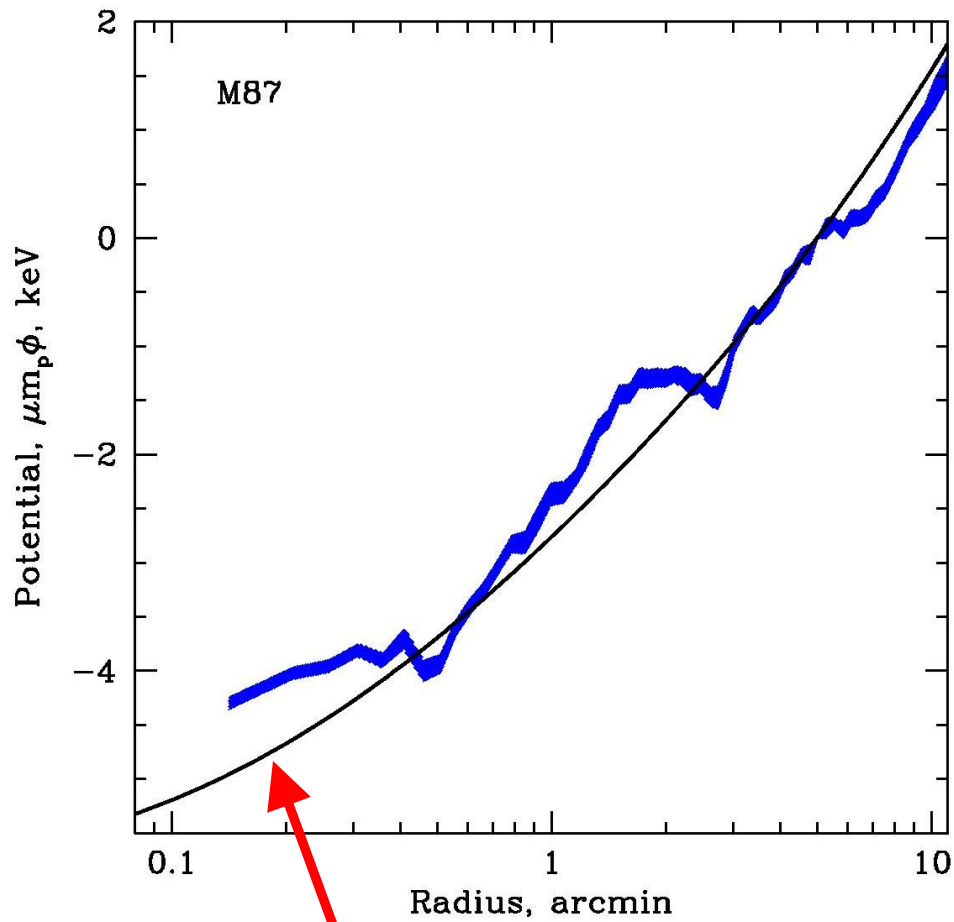
Forman+, 07

# Shock wave in M87

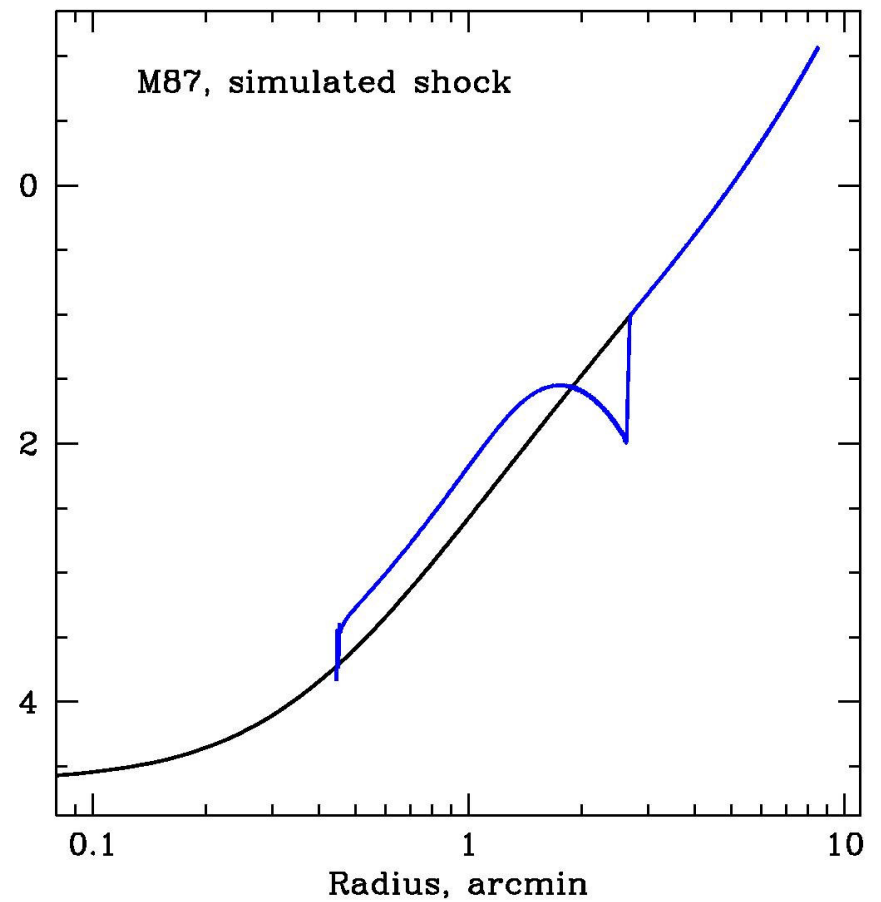
$$E = 5 \cdot 10^{57} \text{ ergs}$$



## Observations

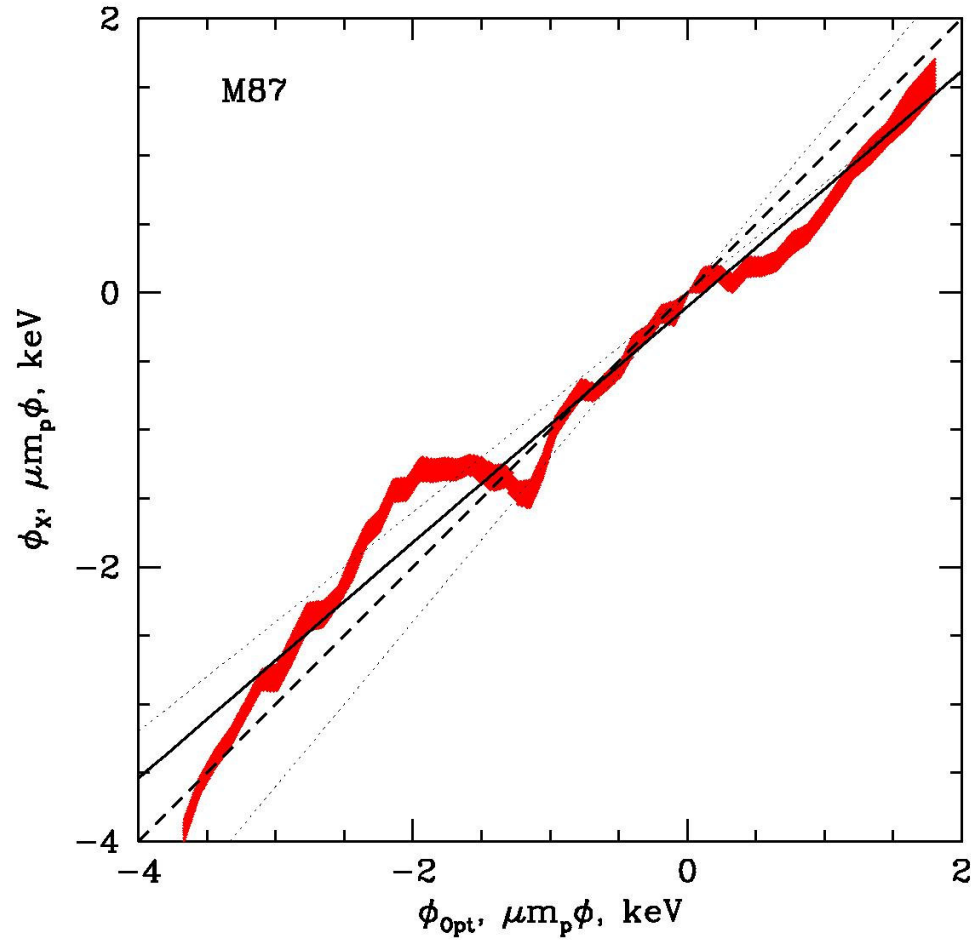


## Model predictions



Potential from stellar kinematics [Romanowsky & Kochanek, 2001]

M87



$$\varphi_X(r) \approx 0.85 \varphi_{opt}(r) + C$$

$$U_{CR} + \frac{H^2}{8\pi} + U_{turb} = 0.15 U_{thermal}$$

# Conclusions

1. Cosmic rays, magnetic fields and microturbulence make ~10% of the total gas pressure in cluster cores  
Mahdavi et al., 2008, Zhang et al. 2008
2. Hydrostatic - (surprisingly) good approximation for ICM (in cores)

$$U_{CR} + \frac{H^2}{8\pi} + U_{turb} \approx 0.1 - 0.2 U_{thermal}$$

GLAST      Faraday      X-ray calorimeters